

# Shadow of Kerr black hole surrounded by an angular Gaussian distributed plasma

Zhenyu Zhang

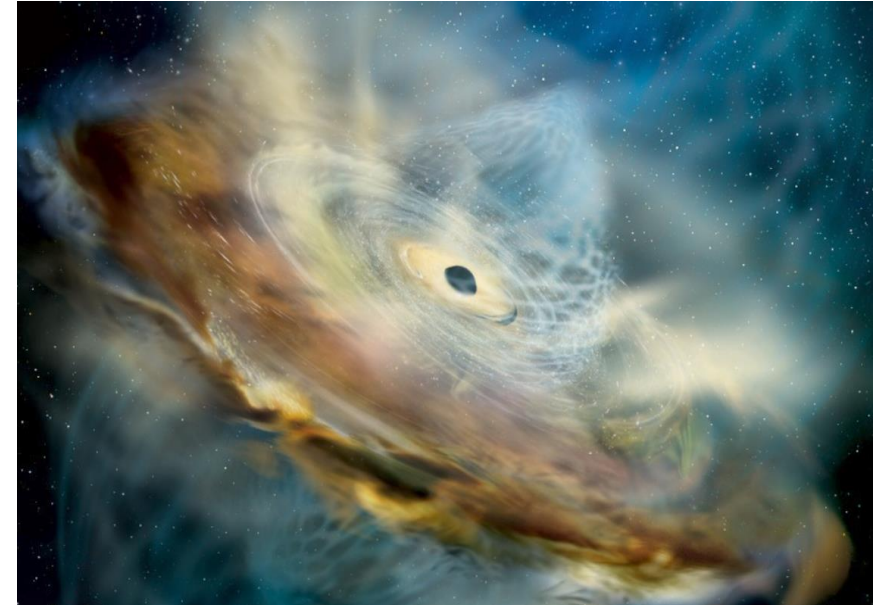
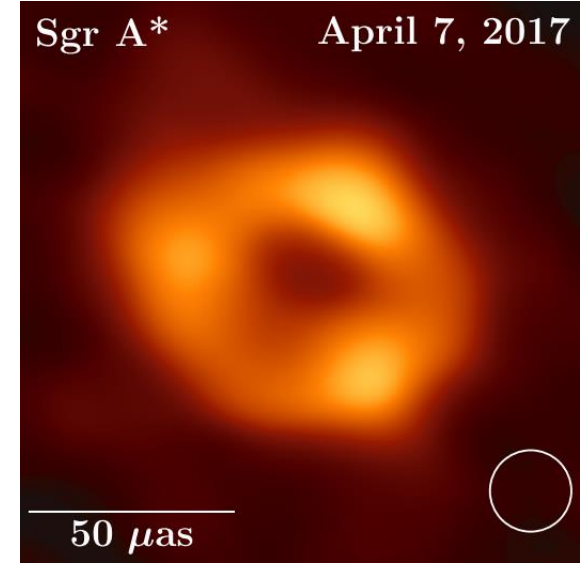
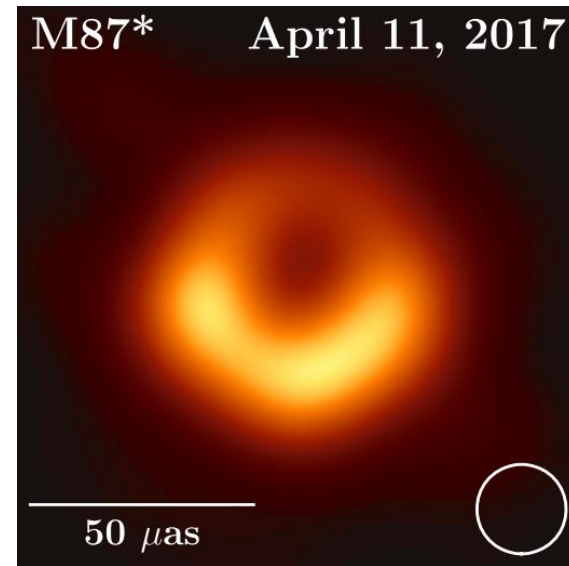
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Based on arXiv: 2206.04430

Collaborators: Haopeng Yan, Minyong Guo and Bin Chen

# Motivation

- Event Horizon Telescope
  - Physical information near black holes
    - Spacetime background
    - Accretion flow
- Accretion Disk
  - Made of **plasma**
  - Type
    - Geometric thin, optical thick
    - **Geometric thick, optical thin**



(Image Source: <https://www.nasa.gov/>)

# Theoretical basis

- Hamiltonian of photon motion in a cold non-magnetized plasma (Synge 1960)

$$H = \frac{1}{2} (g_{\mu\nu} p^\mu p^\nu + \omega_p^2)$$

$\omega_p$  : plasma frequency

- According to plasma physics

$$\omega_p^2(x) = \frac{4\pi e^2}{m_e} N(x) \propto N(x)$$

- Analytical methods are only available for some special cases

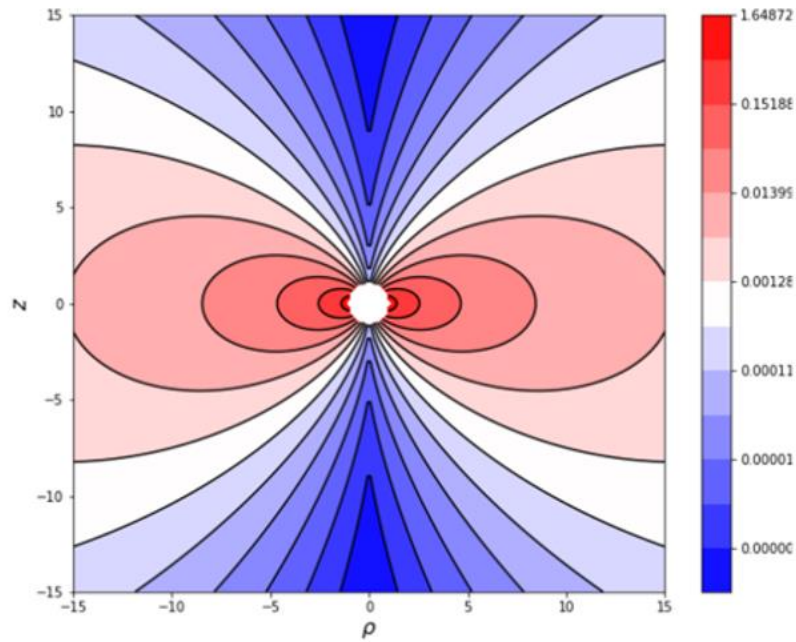
$$\omega_p^2(r, \theta) = \frac{f_r(r) + f_\theta(\theta)}{r^2 + a^2 \cos^2 \theta}$$

- Numerical method is applied in our work for general case

# Plasma distribution models

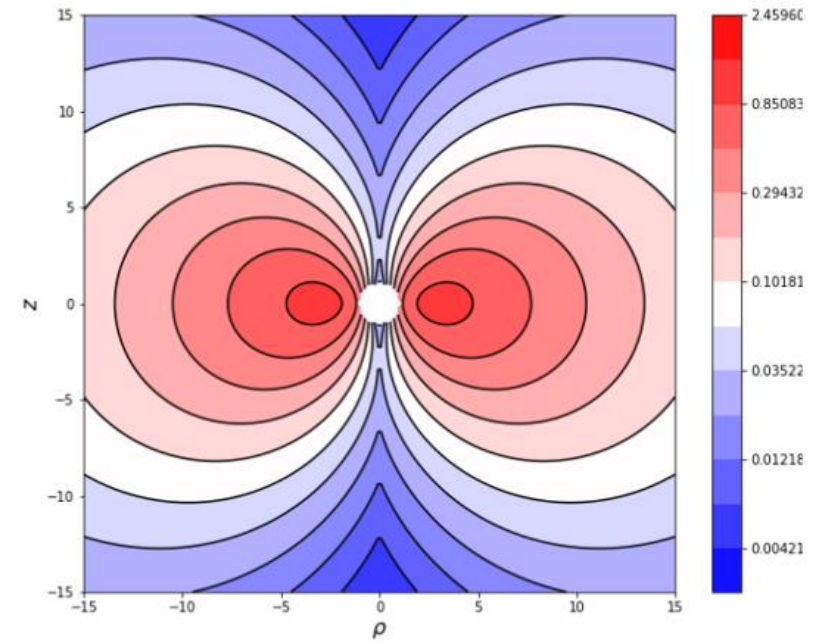
Model A (two parameters)

$$\omega_{pA}^2(r, \theta) = \frac{k_A}{r^2} e^{-\frac{(\theta - \frac{\pi}{2})^2}{2\xi_\theta^2}}$$

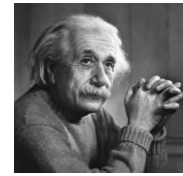
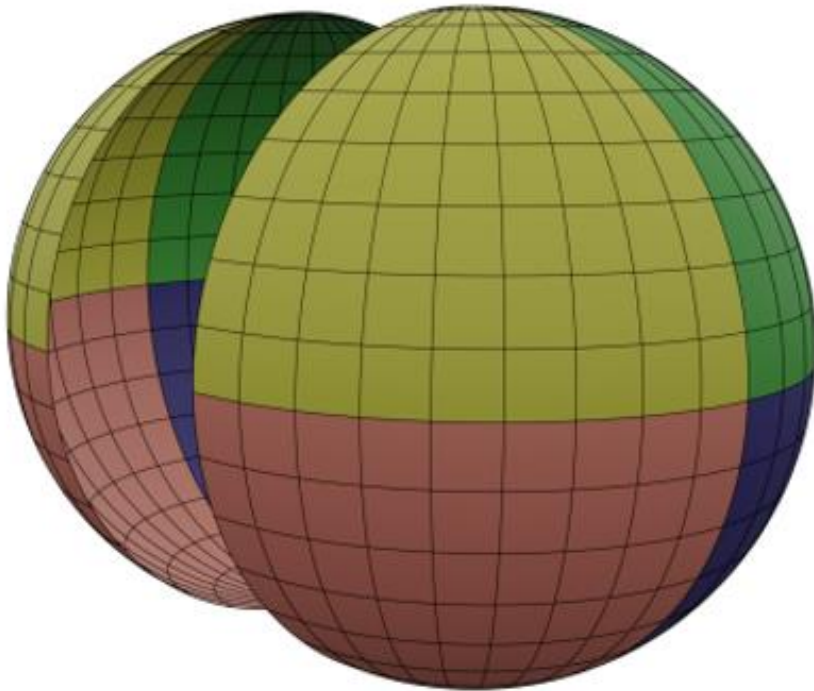


Model B (four parameters)

$$\omega_{pB}^2(r, \theta) = k_B e^{-\frac{(\log \frac{r}{r_m})^2}{2\sigma^2}} e^{-\frac{(\theta - \frac{\pi}{2})^2}{2\xi_\theta^2}}$$



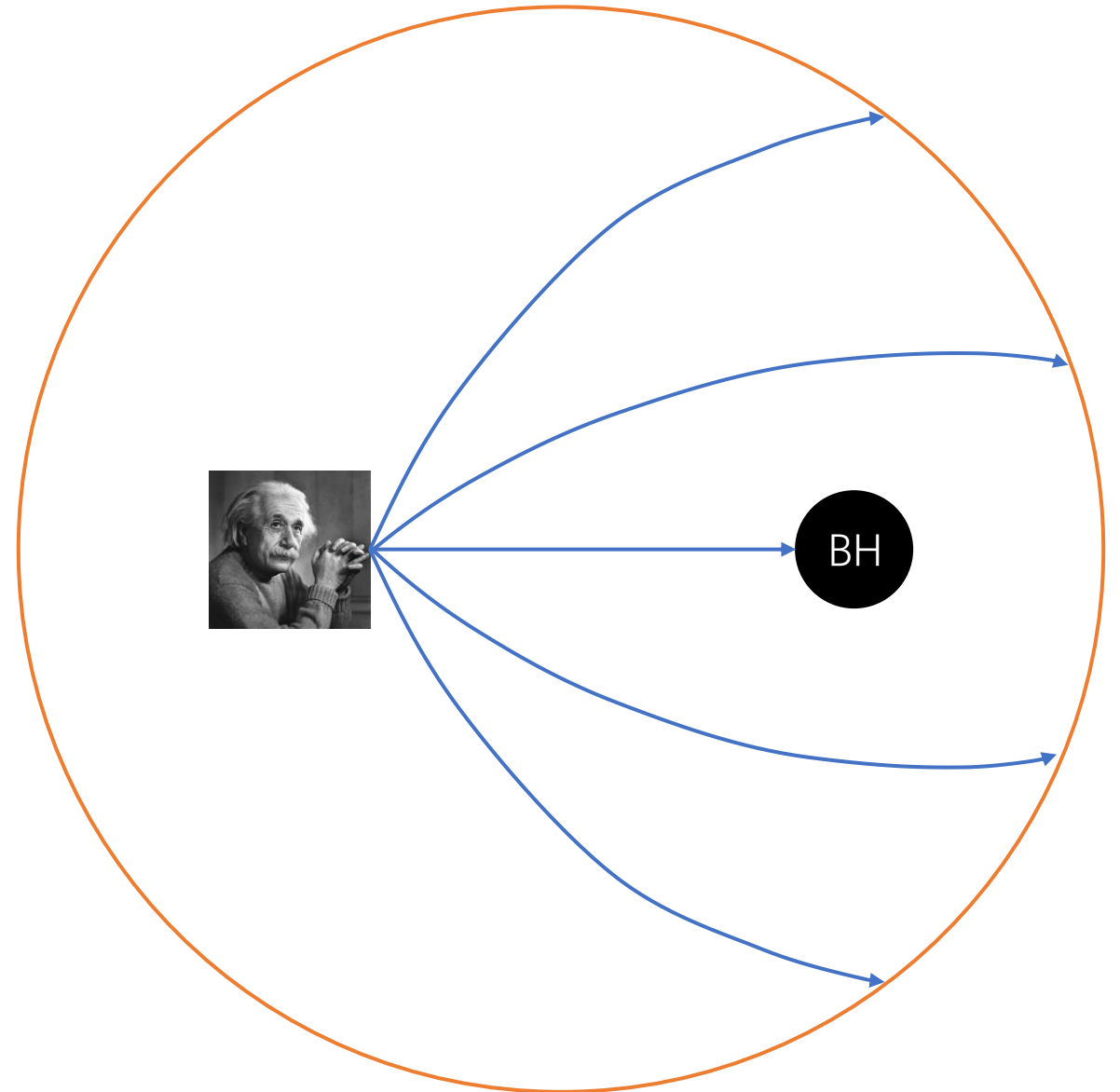
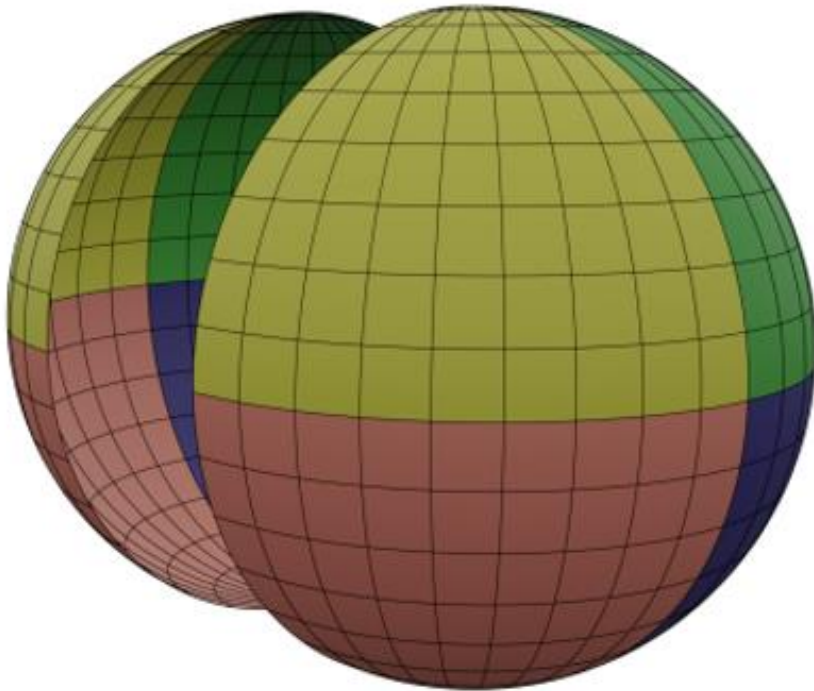
# Numerical method: backward ray-tracing



$$\dot{p}_\mu = -\frac{\partial H}{\partial x^\mu}$$
$$\dot{x}^\mu = \frac{\partial H}{\partial p_\mu}$$



# Numerical method: backward ray-tracing

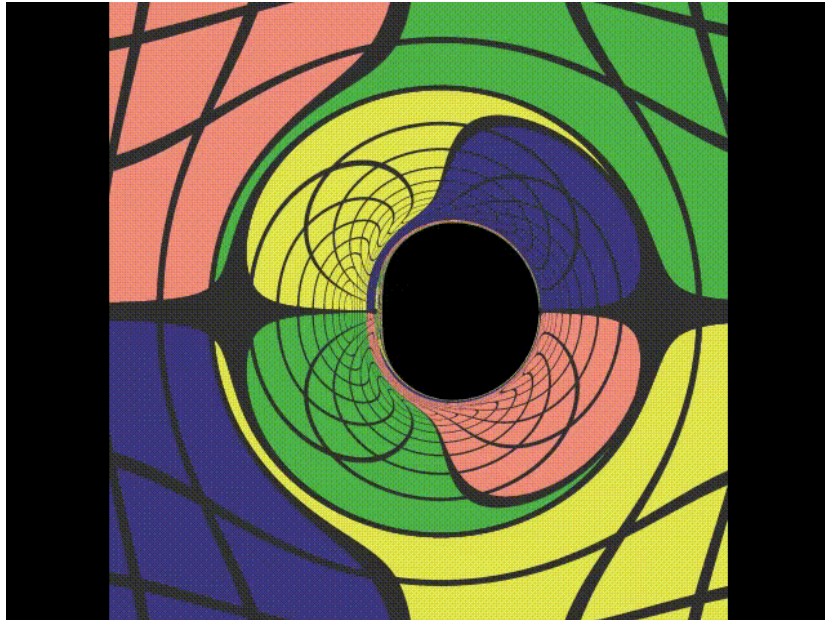




# Results of model A

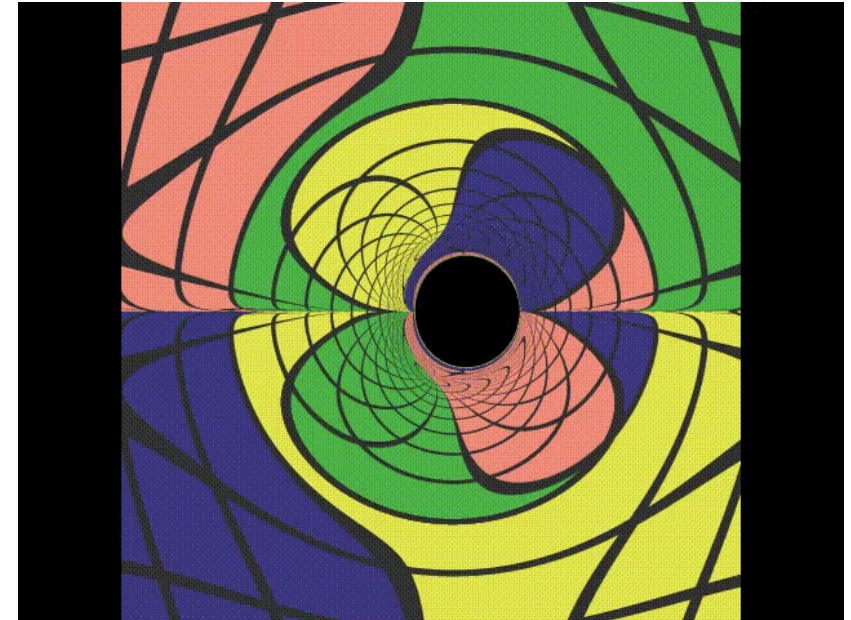
$$\omega_{pA}^2(r, \theta) = \frac{k_A}{r^2} e^{-\frac{(\theta - \frac{\pi}{2})^2}{2\xi_\theta^2}}$$

Change  $k_A$ : 1 ~ 26



$$a = 0.998, \theta_{obs} = \frac{\pi}{2}, \xi_\theta = 0.36$$

Change  $\xi_\theta$ : 0.09 ~ 0.72



$$a = 0.998, \theta_{obs} = \frac{\pi}{2}, k_A = 16$$

Set  $G = c = M = 1$

The effect of changing the position of the observer is also studied in our paper

# Results of model B

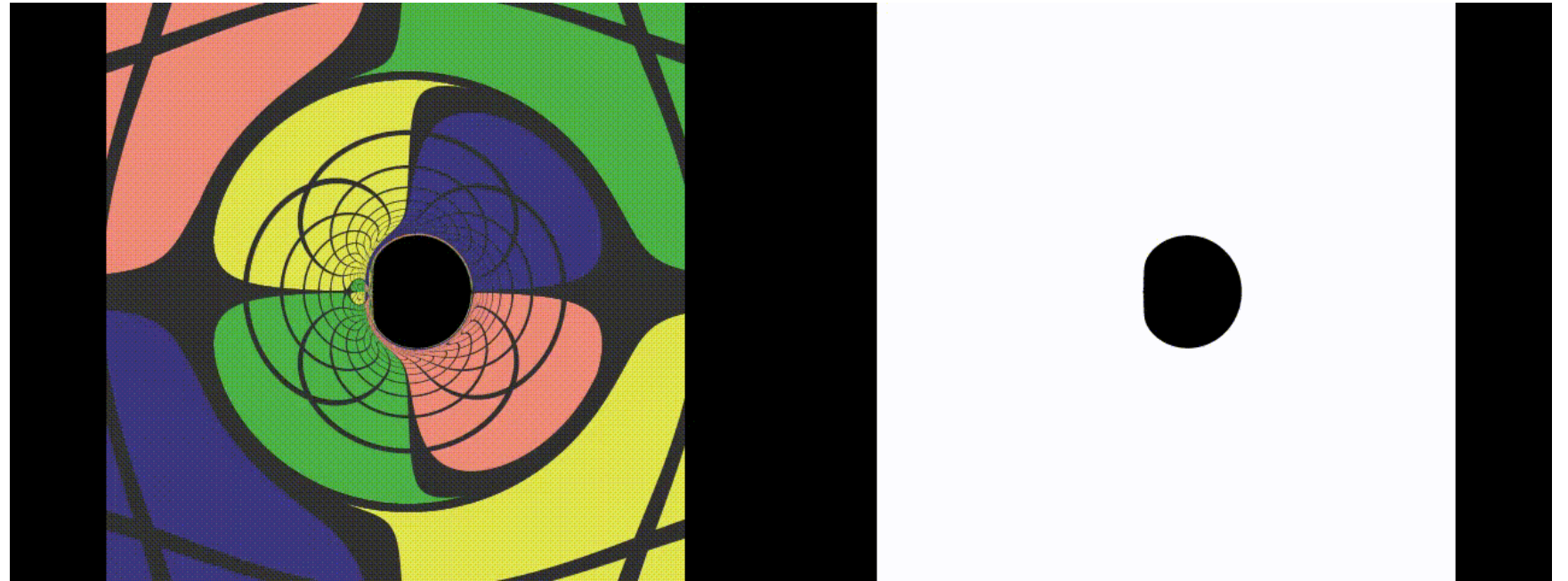
$$\omega_{pB}^2(r, \theta) = k_B e^{-\frac{\left(\log \frac{r}{r_m}\right)^2}{2\sigma^2}} e^{-\frac{\left(\theta - \frac{\pi}{2}\right)^2}{2\xi_\theta^2}}$$

$$r_1 = \frac{r_m - r_h}{10} i, \quad i = 1, 2, \dots, 10$$

→  $(r_m, \sigma) \rightarrow (\mathbf{r_m}, \mathbf{i})$  a smaller  $i$  corresponds to a faster decay rate

$$r_{1,2} = r_m e^{\pm \sqrt{2 \log 10} \sigma}$$

Change  $r_m$ : 2 ~ 12



Set  $G = c = M = 1$

$$a = 0.998, \theta_{obs} = \frac{\pi}{2}, k_B = 0.8, \xi_\theta = 0.36, i = 10$$



# Results of model B

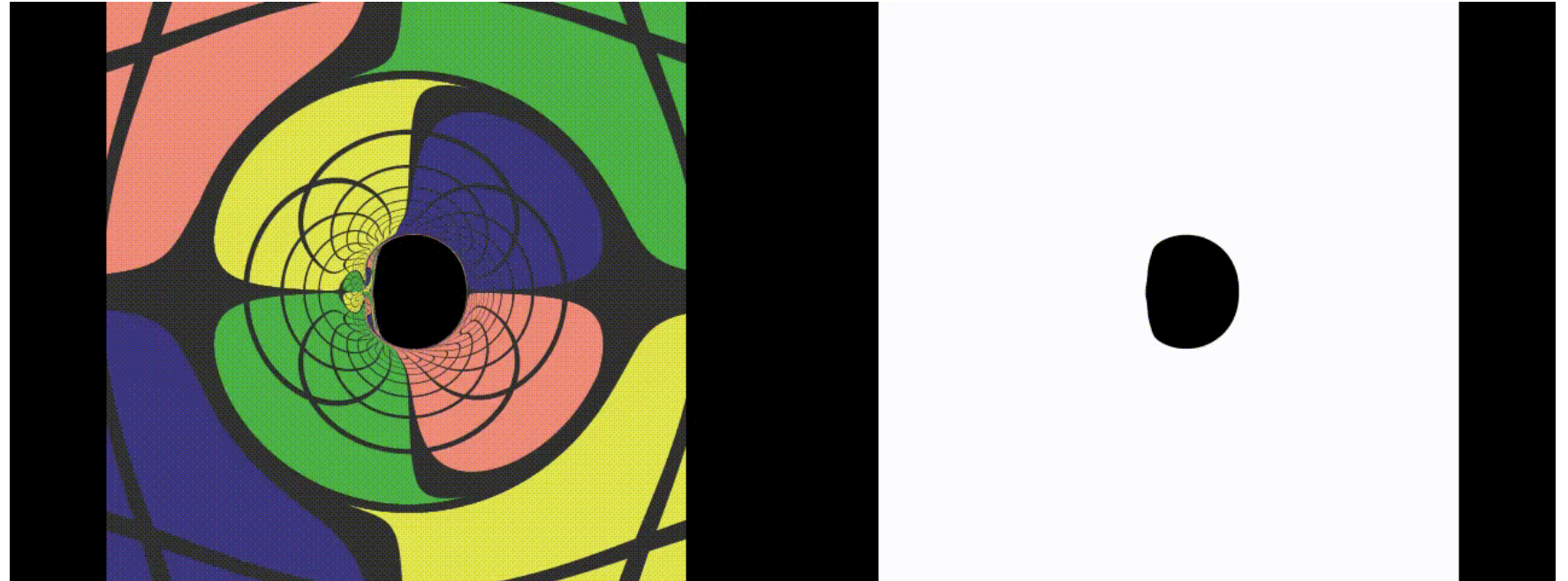
$$\omega_{pB}^2(r, \theta) = k_B e^{-\frac{\left(\log \frac{r}{r_m}\right)^2}{2\sigma^2}} e^{-\frac{\left(\theta - \frac{\pi}{2}\right)^2}{2\xi_\theta^2}}$$

$$r_1 = \frac{r_m - r_h}{10} i, \quad i = 1, 2, \dots, 10$$

→  $(r_m, \sigma) \rightarrow (\mathbf{r_m}, \mathbf{i})$  a smaller  $i$  corresponds to a faster decay rate

$$r_{1,2} = r_m e^{\pm \sqrt{2 \log 10} \sigma}$$

Change  $i$ : 10 ~ 1

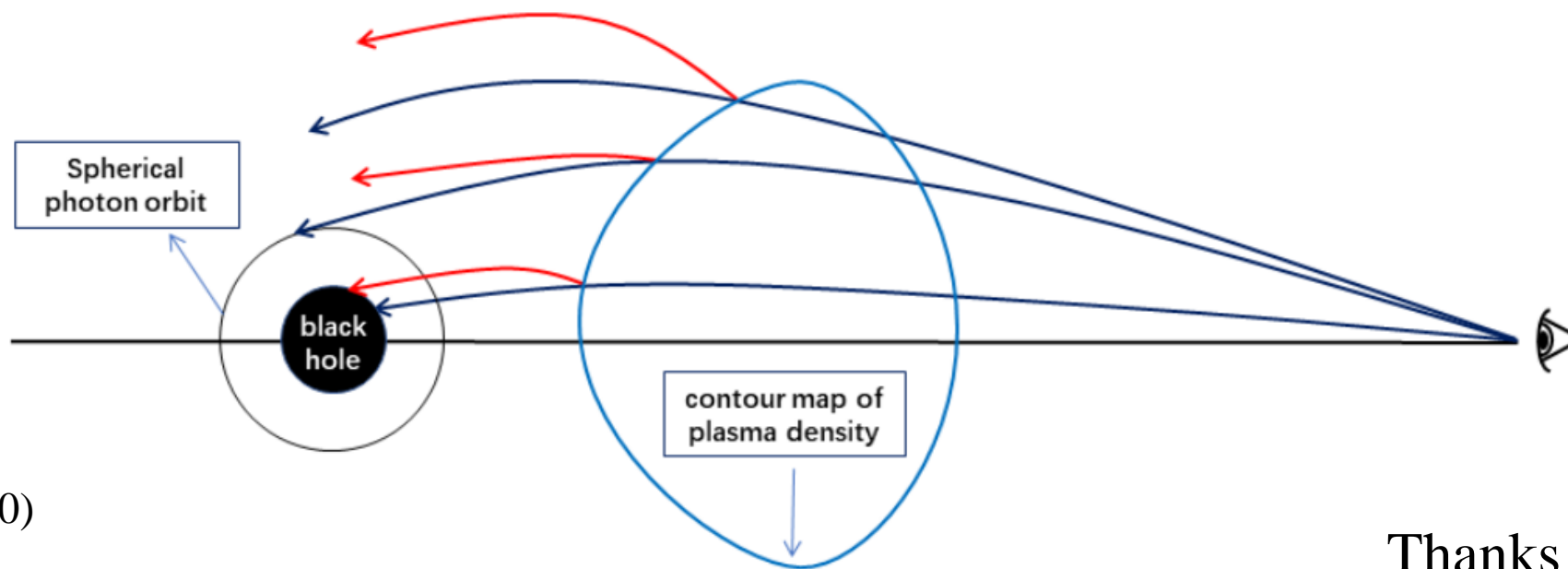
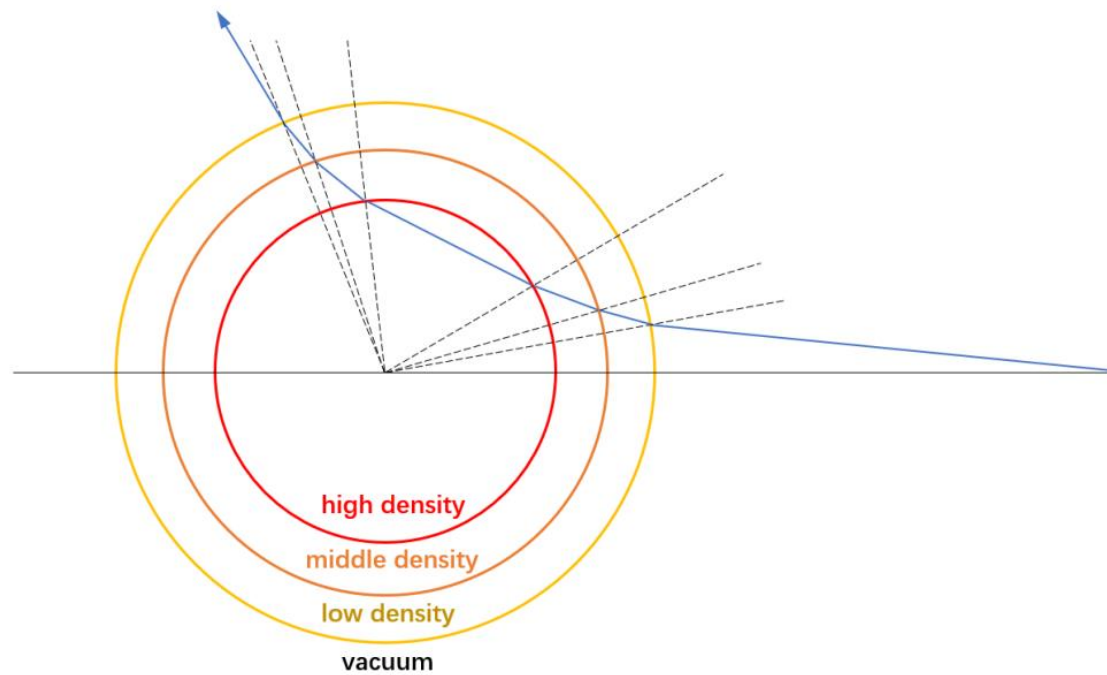


Set  $G = c = M = 1$

$$a = 0.998, \theta_{obs} = \frac{\pi}{2}, \xi_\theta = 0.36, r_m = 3$$

# Qualitative explanation

- The refractive index  $n^2 = 1 - \frac{\omega_p^2}{\omega^2} < 1$
- Fewer rays hit the black hole



See our paper  
(arXiv:2206.04430)  
for more details

Thanks for listening!